# Quantum Chemical and RRKM Investigation of the Elementary Channels of the Reaction $\mathrm{C}_{6} \mathrm{H}_{6}+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ 

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#### Abstract

We performed a computational study of an important reaction in the combustion of hydrocarbons, $\mathrm{C}_{6} \mathrm{H}_{6}$, + O ( ${ }^{3} \mathrm{P}$ ), using ab initio and RRKM methods. Density functional theory (B3LYP) was used to optimize geometries and obtain molecular vibrational frequencies, and complete basis set extrapolation (CBS-QB3) was used to obtain the energies for the reactants several transition states and products. The initial formation of a stabilized adduct is characterized by a barrier of $4.9 \mathrm{kcal} \mathrm{mol}^{-1}$, in good agreement with the measured activation energy for this reaction. All product channels originate from rearrangement or decomposition of this adduct, which our calculations suggest is a triplet ground state. All of our ab initio calculations are thus conducted on the triplet surface. There are several products that are energetically accessible at combustion temperatures, but the formation of phenoxy radical, $\mathrm{C}_{6} \mathrm{H}_{5} \mathrm{O}$, and H atom dominate the product slate at low temperatures. Rearrangement to form formylcyclopentadiene, $\mathrm{C}_{5} \mathrm{H}_{5} \mathrm{CHO}$, is also important at low temperatures, and the decomposition of this species to form cyclopentadienyl radical, $\mathrm{C}_{5} \mathrm{H}_{5}$, and HCO may be significant at higher temperatures. Rearrangement to form phenol is unimportant at all temperatures.


## Introduction

The motivation for this work is to gain an understanding of one of the more important reactions involved in the high temperature oxidation of benzene. The combustion of aromatics in general, and benzene in particular, is important for several reasons. First, aromatic compounds are a large component of many of our fuels; for example, they comprise approximately $32 \%$ of gasoline. ${ }^{1}$ Second, the pyrolysis of aromatic compounds is well-known to have high tendencies toward the production of soot and polycyclic aromatic hydrocarbons (PAH). ${ }^{2}$ Finally, benzene and other aromatics are formed in flames of aliphatic fuel. A detailed understanding of the combustion mechanisms of benzene and other aromatics could help us design better processes to mitigate its adverse effects.

Interest in this study was prompted by an elementary reaction modeling study of low-pressure benzene combustion. ${ }^{3}$ By employing the most commonly assumed product channel for the title reaction (i.e., $\mathrm{C}_{6} \mathrm{H}_{5} \mathrm{O}+\mathrm{H}$ ) and literature values for its rate constant, the concentration of phenoxy radical $\left(\mathrm{C}_{6} \mathrm{H}_{5} \mathrm{O}\right)$ is overpredicted by approximately 2 orders of magnitude. This problem with the phenoxy profile led to two possible hypotheses: (1) the production rate of phenoxy in our reaction mechanism is too high due to either incorrect rate constants or erroneous product channels or (2) the destruction rate of phenoxy in the mechanism is too low either due to errors in the literature rate constants or missing channels. The focus of this work is on the former hypothesis. That is, we report on our computer modeling investigation of the product channels for the reaction of benzene with triplet O atom and the rates for each channel.

[^0]The approach used in this work was ab initio quantum mechanical modeling of the reactants, unimolecular stabilized adducts, and transition states and subsequent RRKM modeling to determine rate constants. The results of our study provide insight into the mechanism for this important reaction.
The title reaction, $\mathrm{C}_{6} \mathrm{H}_{6}+\mathrm{O}\left({ }^{3} \mathrm{P}\right)$, has a large number of energetically feasible product channels. Fortunately, many experimental studies have been carried out ${ }^{4-13}$ with a few of the more pertinent results being summarized as follows. Nicovich et al. ${ }^{6}$ measured absolute rate constants $(A=2.8 \times$ $10^{13} \mathrm{~cm}^{3} \mathrm{~mol}^{-1} \mathrm{~s}^{-1}, E_{\mathrm{a}}=4.9 \mathrm{kcal} \mathrm{mol}^{-1}$ ) for the reaction using flash photolysis-resonance fluorescence over a wide range of temperatures (298-950 K). This group found that their Arrhenius parameters were consistent with most of the previous measurements, which were made over a more narrow temperature span. Furthermore, the data were consistent with addition being the product channel, although the products were not directly measured. The most recent work by Ko et al. ${ }^{7}$ using a method similar to Nicovich et al. over a temperature range of $600-1330 \mathrm{~K}$ returned Arrhenius parameters and conclusions very similar to those of ref 4 . Benzene combustion modelers have employed the addition channel (product $=$ phenol),,${ }^{14,15}$ the abstraction channel (product $=$ phenyl +OH ),,${ }^{14,15}$ and the addition/elimination channel $($ product $=$ phenoxy +H$) .{ }^{3,16,17}$

Reactions of ground-state $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ with hydrocarbons have been extensively studied in the previous few decades. For alkanes, the only reported product channel is the abstraction of hydrogen leading to an alkyl radical and OH . The literature ${ }^{18}$ on this type of reaction is largely in agreement, with $A$-factors near the gas kinetic limit ( $>1.0 \times 10^{14} \mathrm{~cm}^{3} \mathrm{~mol}^{-1} \mathrm{~s}^{-1}$ ) and activation energies of 6 to $10 \mathrm{kcal} \mathrm{mol}^{-1}$, depending upon the strength of the $\mathrm{C}-\mathrm{H}$ bond being broken. The picture becomes dramatically less clear when unsaturates are considered. In this case, the O atom can


Figure 1. Reaction scheme considered for this study. All unimolecular product channels are formed in the triplet state.
still abstract hydrogen, but it can also undergo electrophillic addition to the molecule forming a triplet bi-radical followed by a number of possible chemically activated decomposition channels. For example, in the reaction $\mathrm{O}\left({ }^{3} \mathrm{P}\right)+\mathrm{C}_{2} \mathrm{H}_{4}$ the following reaction product channels are reported: $\mathrm{OH}+\mathrm{C}_{2} \mathrm{H}_{3}$, $\mathrm{H}+\mathrm{CH}_{2} \mathrm{CHO}, \mathrm{H}_{2}+\mathrm{CH}_{2} \mathrm{CO}, \mathrm{H}+\mathrm{CH}_{3} \mathrm{CO}, \mathrm{CH}_{3}+\mathrm{HCO}, \mathrm{CH}_{2}$ $+\mathrm{CH}_{2} \mathrm{O}$, and oxirane. Most O atom reactions involve free radical branching, which further underscores their importance in correctly modeling combustion chemistry.

Several studies involved the analysis of the title reaction using crossed molecular beams. One of the early studies ${ }^{9}$ observed phenol, CO, and a polymeric material as products, but subsequent studies have found only a small amount of CO produced directly by this reaction. Probably the most detailed crossed molecular beeam investigation into possible reaction products was done by Sibener et al.. ${ }^{5}$ These authors observed an addition product (identified as phenol), a hydrogen elimination product (identified as phenoxy radical) and a small amount ( $<5 \%$ ) of CO elimination. They also reached the conclusion that the energy barrier was in the range of $3-5 \mathrm{kcal} \mathrm{mol}^{-1}$. Because the experiment involved nonthermal supersonic beams, this group did not compute Arrhenius parameters.

The most important result of the experimental studies with regard to our benzene modeling work is that they all appear to be in agreement on the rate of the reaction and that addition of O atom dominates over the abstraction channel. The final outcome of the adduct is not directly identified in any of the studies. The point of this study is to determine the product distribution from quantum chemical modeling, and to provide rate constants for these elementary reactions.

The reaction of ground state $\left({ }^{3} \mathrm{P}\right)$ atomic oxygen with benzene is most likely to proceed first by addition to the aromatic ring as in reaction 1 in Figure 1. Abstraction of a hydrogen to form OH and phenyl radicals is unlikely. Nicovitch et al. ${ }^{6}$ found identical rate constants with deuterated benzene $\mathrm{C}_{6} \mathrm{D}_{6}$, leading to the conclusion that the reaction mainly occurs over the addition pathway. One can expect a small barrier for this reaction
since the aromatic delocalization on the ring is broken. Similarly, a small barrier is found in the reaction of hydrogen atom to benzene. ${ }^{19}$ One can easily estimate the heat of this reaction using known bond dissociation energies and estimates of resonance stabilization. The enthalpy of this reaction should be

$$
\begin{aligned}
& \Delta H_{\text {react } 1}=D^{\circ}(\mathrm{C}=\mathrm{C} \rightarrow \mathrm{C}-\mathrm{C})+\Delta_{\text {stab }} H(\text { benzene })- \\
& D^{\circ}(\mathrm{C}-\mathrm{O})-\Delta_{\text {stab }} H(\text { cyclohexadienyl })
\end{aligned}
$$

where $D^{\circ}(\mathrm{C}=\mathrm{C} \rightarrow \mathrm{C}-\mathrm{C})$ is the dissociation energy of a $\pi$-bond in a double bond ( 58 kcal mol-1), ${ }^{20} \Delta_{\text {stab }} H$ (benzene) is the aromatic stabilization energy of benzene $\left(37 \mathrm{kcal} \mathrm{mol}^{-1}\right),{ }^{21} D^{\circ}$ -$(\mathrm{C}-\mathrm{O})$ is a bond dissociation energy of a $\mathrm{C}-\mathrm{O}$ bond ( 82 kcal $\left.\mathrm{mol}^{-1}\right)^{22}$ and $\Delta_{\text {stab }} H$ (cyclohexadienyl) is the stabilization energy of the 2,4 -cyclohexadienyl radical ( $20 \mathrm{kcal} \mathrm{mol}-1$ ). ${ }^{23}$ Using these values, one obtains the value $\Delta H_{\text {react } 1}=-7 \mathrm{kcal} \mathrm{mol}^{-1}$. The unpaired electron on the oxygen atom is separated from the unpaired electron on the ring by a $\mathrm{sp}^{3}$ carbon, and thus, one expects the ground state of $\mathbf{1}$ to be a triplet state. One would expect that the products formed from the reaction of atomic oxygen with benzene would arise from reactions of this chemically activated adduct. There are several possible energetically accessible exit channels that we considered for this study, shown in Figure 1.

## Computational Approach

Our approach was to calculate the energies and vibrational frequencies of the molecules $\mathbf{1 - 1 3}$, benzene, CO and HCO , the atoms H and $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ and the transition states TS1-TS13 using electronic structure methods ${ }^{24}$ and evaluate rate constants for reactions shown in Figure 1, using Rice Ramsperger Kassel Marcus (RRKM) reaction theory. ${ }^{25}$ In this investigation, Gaussian $98^{26}$ was used to calculate molecular properties and our own code was used for the RRKM calculations. Because the oxygen atom is a triplet in its ground state, we have conducted our ab initio calculations on the triplet surface. Our B3LYP and CBS-

TABLE 1: Arrhenius Fits of the Rate Constants for the Disappearance of Benzene and Oxygen, as Compared to Several Experimental Values

| source | $A\left(\mathrm{~cm}^{3} \mathrm{~mol}^{-1} \mathrm{~s}^{-1}\right)$ | $E\left(\mathrm{kcal} \mathrm{mol}^{-1}\right)$ |
| :--- | :---: | :---: |
| Tappe et al. ${ }^{13}$ | $2.1 \times 10^{13}$ | 4.4 |
| Nicovitch et al. ${ }^{6}$ | $2.78 \times 10^{13}$ | 4.9 |
| Ko et al. $^{7}$ | $3.22 \times 10^{13}$ | 5.1 |
| RRKM $^{a}$ | $1.36 \times 10^{14}$ | 6.1 |
| ${ }^{a}$ This work, 100 Torr. |  |  |

QB3 calculations suggest that the adduct $\mathbf{1}$ has a triplet ground state. All of the products, with the exception of the radical products, are singlets in their ground states. We did not attempt to calculate the intersystem crossing rates for hopping from the triplet surface to the singlet surface. The lifetimes of the intermediates in this study are very short at combustion temperatures (see below) compared to typical triplet lifetimes. ${ }^{27}$

Electronic Structure Calculations. Initially, geometries and energies were determined using MP4/6-31G(d)//HF/6-31G(d). However, this method of calculation proved time-consuming and inaccurate as is shown in Table 2 for reaction 1. Higher energy accuracy was obtained using the density functional technique, B3LYP, ${ }^{28}$ and the correlation consistent basis set (ccpVDZ) developed by Dunning. ${ }^{29}$ The B3LYP technique has been shown to produce accurate energies at a low computational cost. ${ }^{30}$ Furthermore, we have found that vibrational frequencies for stable compounds and radicals can be accurately determined using B3LYP. ${ }^{31-35}$

The results of our B3LYP calculations of the transition state for reaction 1, TS1, are gathered in Table 2. Though the results from the B3LYP method are an improvement over HF/6-31G(d) and MP4/6-31G(d)//HF/6-31G(d), the transition state energy is still low relative to experiment. Earlier studies using the B3LYP technique have also demonstrated that experimental activation energies are often underestimated. ${ }^{36-49}$ Consistent with this, our calculated barrier was $4-5 \mathrm{kcal} \mathrm{mol}^{-1}$ lower than the experimental value. This error did not improve with a larger basis set (cc-pVQZ). On the other hand calculations using B3LYP-optimized geometries and CCSD and QCISD energies are about $5 \mathrm{kcal} \mathrm{mol}^{-1}$ over experimental values. CBS-Q, which optimizes geometries with MP2, is also approximately 5 kcal $\mathrm{mol}^{-1}$ too high. To improve our energy calculations, we used a newly developed CBS extrapolation technique, CBS-QB3. ${ }^{50}$ This technique obtains molecular geometries and vibrational frequencies using B3LYP. Using the improved geometries and vibrational frequencies of B3LYP relative to MP2, the CBSQB3 technique has been shown to produce energies accurate to approximately $2.8 \mathrm{kcal} \mathrm{mol}^{-1} .{ }^{50}$ Our calculated energy barrier using CBS-QB3 is in excellent agreement with the experimental value as shown in Table 1.

The location of transition states in this study was determined using the Synchronous Transit Quasi Newton (STQN) method ${ }^{51}$ available in Gaussian 98. In most cases, this technique was successful at locating a saddle point that appeared to connect the reactant(s) with the product(s). When this was unsuccessful, we would conduct a potential energy surface (PES) calculation, where the reaction coordinate was varied systematically and all other degrees of freedom were allow to relax. This technique provided a rough estimate of the geometry and energy of the transition state. The actual transition state could then be located easily using a Berny minimization. TS6, TS10, and TS11 were determined in this manner. There were some transition states where even this technique failed, and we could only get rough
estimate of the transition state from the PES calculations. This was the case for reactions 8 and 9 , which have large geometry changes in the transition state and high barriers. Thus, these channels do not contribute significantly.

RRKM Kinetics. From RRKM theory, the microcanonical rate constant is derived from the density of states of the unimolecular reactant and sum of states of the transition state. The microcanonical rate constant, in the notation of Holbrook, ${ }^{52}$ is

$$
\begin{equation*}
k(E)=L^{\ddagger} \frac{Q^{\dagger}}{Q} \frac{W\left(E^{\dagger}\right)}{h \rho(E)} \tag{1}
\end{equation*}
$$

The impact of molecular properties on the RRKM theory can easily be seen in eq 1 , where $L^{\ddagger}$ is the statistical reaction path degeneracy, $Q$ is the partition function, $W(E)$ is the sum of states of the transition state, $h$ is Planck's constant, and $\rho(E)$ is the density of states. The superscript " $\dagger$ " refers to the transition state properties. The sum of states is the integral of the density of states of the transition state structure from its zero point up to the energy of evaluation. The density of states of a molecule was calculated by assuming that all modes (vibrations, rotations, hindered rotations) are separable, and each vibration can be modeled as a harmonic oscillation. Each mode is convoluted together with the Stein-Rabinovitch ${ }^{53}$ algorithm to get all the possible quantum energy levels of the molecule. RRKM theory is applicable to systems that adhere to the assumption of ergodicity; that is, the energy imparted in a collision is randomly distributed throughout all available modes of motion on the time scale of reaction.

As can be found in reference 52 the overall rate constant is determined by assuming a pseudo steady state on all energetically activated species, and integrating the resulting microcanonical expression over energy ranging from the barrier height to infinity (or a high enough energy that the Boltzmann distribution function annuls the contribution at that energy level).

For the formation of stabilized adduct, the rate constant is

$$
\begin{align*}
& k_{\text {stab }, I}=K_{\mathrm{eq}} \int \\
& 1+\frac{k_{-1}}{\beta \omega}+\frac{k_{2}}{\beta \omega}+\frac{k_{4}}{\beta \omega}-\frac{\frac{k_{3}}{\beta \omega} \frac{k_{-3}}{\beta^{\prime} \omega}}{1+\frac{k_{-3}}{\beta^{\prime} \omega}+\frac{k_{10}}{\beta^{\prime} \omega}+\frac{k_{11}}{\beta^{\prime} \omega}}-\frac{\frac{k_{5}}{\beta \omega} \frac{k_{-5}}{\beta^{\prime \prime} \omega}}{1+\frac{k_{-5}}{\beta^{\prime \prime} \omega}}-\frac{\frac{k_{12}}{\beta^{\prime \prime} \omega} \frac{k_{-12}}{\beta^{\prime \prime \prime} \omega}}{1+\frac{k_{-12}}{\beta^{\prime \prime \prime} \omega}}-\frac{\frac{k_{7}}{\beta \omega}}{\frac{k_{-7}}{\beta^{\prime \prime \prime} \omega}}  \tag{2}\\
& 1+\frac{k_{-7}}{\beta^{\prime \prime \prime \prime} \omega}
\end{align*}
$$

where $P(E)$ is the Boltzmann distribution function divided by the partition function of all active modes in the adduct $\mathbf{1}$. The collision factor $\beta$, evaluated below, is unique to each individual unimolecular stabilization channel. The equilibrium rate constant $K_{\text {eq }}$ was calculated from statistical mechanics, via a ratio of the partition functions of benzene and oxygen reactants and product $\mathbf{1}$, as shown in the equation

$$
\begin{equation*}
K_{\mathrm{eq}}=\frac{Q_{\mathrm{AB}}}{Q_{\mathrm{A}} Q_{\mathrm{B}}} \exp \left(\frac{-\Delta H}{k_{\mathrm{B}} T}\right) \tag{3}
\end{equation*}
$$

Where the partition functions are calculated according to their ground-state energy, and $\Delta H$ is the change in enthalpy between the reactants ( A and B ) and product ( AB ).

TABLE 2: Comparison of Various Computational Methods for the Calculated Energy for Molecules in the Entrance Channel (Reaction 1)

| method | energy in Hartrees (ZPE included) |  |  | barrierheight$\left(\mathrm{kcal} \mathrm{mol}^{-1}\right)$ |
| :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ | $\mathrm{C}_{6} \mathrm{H}_{6}$ | TS1 |  |
| HF/6-31G(d) | -74.783 93 | -230.70314 | -305.486 94 | $-1.5$ |
| MP4/6-31G(d)//HF/6-31G(d) | -74.895 97 | -231.58034 | -306.455 47 | 11.5 |
| B3LYP/cc-pVDZ | -75.068 49 | -232.26296 | -307.329 08 | 1.4 |
| B3LYP/cc-pVQZ | -75.098 19 | -232.34927 | -307.44673 | 0.4 |
| QCISD-T/6-31G(d)//B3LYP/cc-pVDZ | -74.895 90 | -231.49639 | -306.373 17 | 11.9 |
| QCISD(T)/6-31G(d)//B3LYP/cc-pVDZ | -74.896 68 | -231.53150 | -306.410 27 | 11.1 |
| QCISD(T)/3-21G(d)/B3LYP/cc-pVDZ | -74.449 71 | -230.00191 | -304.427 69 | 14.9 |
| B3LYP/6-311G(d,p) | -75.085 39 | -232.30855 | -307.39177 | 1.1 |
| CBS-QB3 | -74.987 64 | -231.88985 | -306.869 11 | 4.9 |
| experimental activation energy ${ }^{6}$ |  |  |  | 4.9 |

For formation of phenoxy radical, the resulting rate expression is

$$
\begin{align*}
& k_{\mathrm{e} \text { lim,2 }}=K_{\mathrm{eq}} \int \\
& 1+\frac{k_{-1}}{\beta \omega}+\frac{k_{2}}{\beta \omega}+\frac{k_{4}}{\beta \omega}-\frac{\frac{k_{3}}{\beta \omega} \frac{k_{-3}}{\beta^{\prime} \omega}}{1+\frac{k_{-3}}{\beta \omega}+\frac{k_{10}}{\beta^{\prime} \omega}+\frac{k_{11}}{\beta^{\prime} \omega}}-\frac{\frac{k_{5}}{\beta \omega} \frac{k_{-5}}{\beta^{\prime \prime} \omega}}{1+\frac{k_{-5}}{\beta^{\prime \prime} \omega}}-\frac{\frac{k_{12}}{\beta^{\prime \prime} \omega} \frac{k_{-12}}{\beta^{\prime \prime \prime} \omega}}{1+\frac{k_{-12}}{\beta^{\prime \prime \prime} \omega}}-\frac{\frac{k_{7}}{\beta \omega} \frac{k_{-7}}{\beta^{\prime \prime \prime} \omega}}{1+\frac{k_{-7}}{\beta^{\prime \prime \prime \prime} \omega}} \tag{4}
\end{align*}
$$

The expression of the rate constant for the formation of stable formylcyclopentadiene is

$$
\begin{align*}
& k_{\text {stab, } 3}=K_{\text {eq }} \int \\
& \frac{k_{-1} \frac{\frac{k_{3}}{\beta \omega}}{1+\frac{k_{-3}}{\beta^{\prime} \omega}+\frac{k_{10}}{\beta^{\prime} \omega}+\frac{k_{11}}{\beta^{\prime} \omega}} P(E) \mathrm{d} E}{1+\frac{k_{-1}}{\beta \omega}+\frac{k_{2}}{\beta \omega}+\frac{k_{4}}{\beta \omega}-\frac{\frac{k_{-3}}{\beta \omega} \frac{k_{5}}{\beta^{\prime} \omega}}{1+\frac{k_{-3}}{\beta^{\prime} \omega}+\frac{k_{10}}{\beta^{\prime} \omega}+\frac{k_{11}}{\beta^{\prime} \omega}}-\frac{\frac{k_{-5}}{\beta^{\prime \prime} \omega}}{1+\frac{k_{-5}}{\beta^{\prime \prime} \omega}}-\frac{\frac{k_{12}}{\beta^{\prime \prime} \omega} \frac{k_{-12}}{\beta^{\prime \prime \prime} \omega}}{1+\frac{k_{-12}}{\beta^{\prime \prime \prime} \omega}}-\frac{\frac{k_{7}}{\beta \omega} \frac{k_{-7}}{1+\frac{k_{-7}^{\prime \prime \prime} \omega}{\beta^{\prime \prime \prime \prime} \omega}}}{15}} \tag{5}
\end{align*}
$$

The formation of the cyclopentyl radical and the formyl radical is

$$
\begin{align*}
& k_{\mathrm{e} \text { lim,l0 }}=K_{\mathrm{eq}} \int \\
& 1+\frac{k_{-1}}{\beta \omega}+\frac{k_{2}}{\beta \omega}+\frac{k_{3} k_{10}}{\beta \omega \beta^{\prime} \omega}  \tag{6}\\
& k_{-1} \frac{k_{4}}{1+\frac{k_{-3}}{\beta^{\prime} \omega}+\frac{k_{10}}{\beta^{\prime} \omega}+\frac{k_{11}}{\beta^{\prime} \omega} P(E) \mathrm{d} E} \\
& 1+\frac{\frac{k_{3}}{\beta \omega} \frac{k_{-3}}{\beta^{\prime} \omega}}{\beta^{\prime} \omega}+\frac{k_{10}}{\beta^{\prime} \omega}+\frac{k_{11}}{\beta^{\prime} \omega}-\frac{k_{-5}}{\beta \omega} \frac{k_{-5}}{\beta^{\prime \prime} \omega} \\
& 1+\frac{k_{-5}}{\beta^{\prime \prime} \omega} \\
& \frac{k_{12}}{\beta^{\prime \prime} \omega} \frac{k_{-12}}{1+\frac{k_{-12}^{\prime \prime \prime} \omega}{\beta^{\prime \prime \prime} \omega}}-\frac{\frac{k_{7}}{\beta \omega}}{1+\frac{k_{-7}}{\beta^{\prime \prime \prime} \omega}} \\
& 1+\frac{k_{-7}}{\beta^{\prime \prime \prime} \omega}
\end{align*}
$$

The rates of other channels are computed using a similar logic.
Factors for the reactions $6,8,9$, and 13 were not considered because the zero point energy of the transition states were high, and consequently, the contributions of these reactions would be expected to be negligible.

To account for collisional stability, we follow a method utilizing Gilbert's ${ }^{54}$ modification to the collision factor, which is a correction to the strong collision theory in which every collision completely deactivates a molecule. We also used an empirical collision integral found in Gilbert and Smith ${ }^{55}$ which describes the frequency of collisions based loosely on kinetic
theory. The collision factor, $\beta$, is calculated from an approximation of all the collisions that will activate or deactivate an excited single molecule. The equation for the collision factor follows.

$$
\begin{gather*}
\beta \approx \frac{\left(\frac{\left\langle E_{\mathrm{down}}\right\rangle}{\left\langle E_{\mathrm{down}}\right\rangle+\mathrm{F}_{\mathrm{e}} k_{\mathrm{b}} T}\right)^{2}}{\int_{0}^{E_{\mathrm{o}}} \rho(E) e^{-} E / k^{\mathrm{b}} T\left(1-\frac{F_{\mathrm{e}}}{1-F_{\mathrm{e}}} e^{\mathrm{E}_{\mathrm{o}}}-E / F^{\mathrm{e}}\right) \mathrm{d} E}  \tag{7}\\
F_{\mathrm{e}}=\frac{\int_{E_{\mathrm{o}}}^{\infty} \rho(E) e^{-} E / k^{\mathrm{b}} T \mathrm{~d} E}{\rho\left(E_{\mathrm{o}}\right) e^{-E / k_{\mathrm{b}} T} k_{\mathrm{b}} T} \tag{8}
\end{gather*}
$$

Wang ${ }^{56}$ verifies this formulation of the weak collision factor. The average downward energy per collision $\left\langle E_{\text {down }}\right\rangle$ quantity is taken from a small collection compiled by Wang. The code used to calculate the kinetic rate constants uses the method of Wang, except for some minor additions. These additions include a quantum evaluation of the density of states of hindered internal rotation.

There was only one internal rotation in the mechanism, occurring on the formylcyclopentadiene 3 species, but the barrier to rotation was small (as estimated from the vibrational frequency and internal moment of inertia) enough that approximating the mode as a free rotation did not alter the results.

## Results

As discussed above, RRKM calculations require the energies of reactants and transition states as well as vibrational frequencies and rotational constants. (There are no hindered rotors of importance in this study.) We have collected the values needed for RRKM calculations in Table 3, which shows the calculated energies, including zero point energies, of the product channels and transition states calculated using CBS-QB3. In addition to the absolute energies at 0 K , Table 3 shows the energies (in $\mathrm{kcal} \mathrm{mol}^{-1}$ ) for all species relative to the reactants: $\mathrm{C}_{6} \mathrm{H}_{6}+\mathrm{O}$ $\left({ }^{3} \mathrm{P}\right)$. Vibrational frequencies and rotational constants are also reported in Table 3, and the imaginary frequencies of the transition states are indicated by a succeeding $i$. These molecular parameters are determined at the B3LYP/6-311 g(d,p) level. Figure 2 shows the atom numbering scheme for this study, and Table 4 shows some of the important bond lengths for the transition states, whereas Figure 3 shows the molecular structures of the calculated transitions states.

Figure 4 shows the computed surface for the reactions of the $\mathrm{C}_{6} \mathrm{H}_{6}+\mathrm{O}$ adduct on the triplet surface. The first step of the reaction is the addition of $\mathrm{O}\left({ }^{3} \mathrm{P}\right)$ to benzene, which is exothermic

TABLE 3: Vibration Frequencies,Energies, and Rotational Constants Calculated with CBS-QB3


TABLE 3 (Continued)

| species | energy ${ }^{a}$ (hartree) | $\Delta E^{b}$ rel $\left(\mathrm{kcal} \mathrm{mol}^{-1}\right)$ | $\nu^{c}$ unscaled, A, B, C ( $\left.\mathrm{cm}^{-1}, \mathrm{GHz}\right)$ |
| :---: | :---: | :---: | :---: |
| reaction $7 \quad \mathrm{C}_{6} \mathrm{H}_{6} \mathrm{O} \rightarrow \mathrm{TS} 7 \rightarrow \mathrm{C}_{6} \mathrm{H}_{6} \mathrm{O}(2,4$-cyclohexadienone) |  |  |  |
| 7, 2,4-cyclohexadienone ( $\mathrm{a}^{3}{ }^{3} \mathrm{~A}^{\prime}$ ) | $\begin{array}{cc} \text { reaction } 7 & \mathrm{C}_{6} \mathrm{H}_{6} \mathrm{O} \rightarrow \\ -306.82958 & \mathrm{TS} 7 \rightarrow \mathrm{C}_{6} \mathrm{H}_{6} \mathrm{O} \\ -32.2 \end{array}$ |  | 35, 182, 421, 440, 468, 498, 552, 576, 695, 755, 805, |
|  |  |  | 887, 915, 951, 961, 993, 1101, 1166, 1166, 1261, 1332, |
|  |  |  | $\begin{aligned} & 1389,1406,1429,1442,1511,1594,2963,2966,3163, \\ & 3178,3188,3201 \end{aligned}$ |
|  |  |  | $A, B, C=5.19117,2.64189,1.76944$ |
|  | reaction $8 \quad \mathrm{C}_{6} \mathrm{H}_{6} \mathrm{O} \rightarrow \mathrm{TS8} \rightarrow \mathrm{C}_{6} \mathrm{H}_{6} \mathrm{O}(2,5-$ cyclohexadienone $)$ |  |  |
| TS8 ${ }^{\text {d }}$ | -307.26311 [B3LY | [45.5] |  |
| 8,2,5-cyclohexadienone (ã, ${ }^{3} \mathrm{~A}^{\prime \prime}$ ) | -306.802 30 | -15.0 | $\begin{aligned} & 103,268,347,391,483,505,600,656,694,775,880,917, \\ & 921,950,958,966,1117,1183,1188,1193,1267,1340,1393, \\ & 1411,1456,1513,1598,2898,2917,3164,3164,3192,3193 \\ & A, B, C=5.32290,2.61325,1.77136 \end{aligned}$ |



 $858,902,924,942,975,1051,1122,1205,1269,1383,1429$, $1467,1500,1551,1873,3129,3136,3147,3154,3172,3239$
$\begin{array}{llll}\left.13, n \text {-butadienylketene ( } \tilde{\mathrm{a}}^{3}{ }^{3} \mathrm{~A}^{\prime}\right) & -306.79621 & -12.4 & A, B, C=4.287601,2.417610,1.675447 \\ 29,115,163,259,317,359,435,459,544,631,675,690,\end{array}$ $794,808,888,995,1007,1048,1127,1174,1234,1372,1385$, $1430,1444,1515,2185,3095,3135,3147,3149,3162,3241$
$A, B, C=4.08871,1.85393,1.51318$
${ }^{a}$ CBS-QB3 energies containing ZPE (scaling factor for frequencies is 0.99). ${ }^{b}$ Energies containing ZPE relative to reactants. ${ }^{c}$ Vibrational frequencies and rotational constants are determined at the B3LYP/6-311 g(d,p) level. ${ }^{d}$ Approximate energy from PES calculations. Actual transition states could not be found (see text). Energies are from B3LYP/6-311G(d,p).
by $-8.8 \mathrm{kcal} \mathrm{mol}^{-1}$ and has a barrier of approximately 4.9 kcal $\mathrm{mol}^{-1}$. The chemically activated adduct can react on the triplet surface to form a number of products. The two lowest barriers are formation of phenoxy radical ( $-3.4 \mathrm{kcal} \mathrm{mol}^{-1}$ relative to reactants) and to formylcyclopentadiene ( $-1.8 \mathrm{kcal} \mathrm{mol}^{-1}$
relative to reactants). Collisional stabilization of the formylcyclopentadiene can occur, or the chemically activated aldehyde can go through an elimination reaction to form a cyclopentadienyl and formyl radicals. Using the results from Table 3, we can calculate Transition State Theory rate constants for the


Figure 2. Atom labels for this study.
reactions shown in Figure 1 (Table 5) using the following estimate for the preexponential $A$ factor

$$
\begin{equation*}
A=\frac{k T}{h}\left(\frac{1}{R T}\right)^{\Delta n} \exp (1-\Delta n) \exp \left(\frac{\Delta S^{\dot{\dagger}}}{R}\right) \tag{9}
\end{equation*}
$$

where $\Delta n$ is the change in the number of molecules in going from the reactants to the TS and $\Delta S^{\ddagger}$ is the change in entropy. Entropies are calculated using the harmonic oscillator (scaling factor $=0.97$ ) and rigid rotor approximations with molecular parameters taken from B3LYP calculations. Activation energies are taken from CBS-QB3 calculations. We have used these Arrhenius parameters to estimate the rate constants for reactions at 298 and 1800 K , and these values are shown in Table 3. As these TST calculations predict, the most important high-pressure exit channel for the adduct on the triplet surface is the formation of phenoxy radical, 2, followed by formylcyclopentadiene, $\mathbf{3}$, and benzene oxide, 5 .

Formation of Adduct (Reaction 1). The initial step in this reaction is the addition of oxygen atom onto the aromatic ring. Because the ground state of O is ${ }^{3} \mathrm{P}$, one would expect spin conservation to dictate that the product of this reaction is a triplet adduct. Our calculations on the triplet surface predict that the enthalpy of the triplet adduct is $8.7 \mathrm{kcal} \mathrm{mol}^{-1}$ lower in energy than the reactants in good agreement with the bond energy estimation shown above. Our calculations indicate that the singlet state of the adduct is higher in energy than the triplet though the exact amount was difficult to determine because of spin contamination problems. More extensive multireference calculations are needed to obtain accurate estimates of this splitting. As can be seen by the estimated rate constants in Table 5, the lifetime of the adduct at combustion temperatures is so small that intersystem crossing to the singlet potential energy surface may be minor.

Our calculations show that the potential energy barrier for this reaction step has the largest effect upon the overall activation energy for $\mathrm{O}\left({ }^{3} \mathrm{P}\right)+\mathrm{C}_{6} \mathrm{H}_{6}$. In Table 2, we make a comparison of the calculated barrier height for this step using different computational techniques. As can be seen, CBS-QB3 provides the most reasonable barrier energy ( $4.9 \mathrm{kcal} \mathrm{mol}^{-1}$ ), and the density functional techniques all predict a barrier that is 3-4 $\mathrm{kcal} \mathrm{mol}^{-1}$ too low. As mentioned above, this result is typical for B3LYP. ${ }^{36-49}$

The scheme below shows the molecular orbital structure for this step of the reaction. The reactants, transition state and adduct are all of $\mathrm{C}_{s}$ symmetry. Allowing the reaction to occur without symmetry produced no difference in the results. The transition

TABLE 4: Selected Bond Lengths and Angles for Transition States at B3LYP/6-311G(d,p)

| molecule | bond length ( $\AA$ ) |  | angles (degrees) |  |
| :---: | :---: | :---: | :---: | :---: |
| Benzene ${ }^{a}$ | $\mathrm{C}-\mathrm{H}$ | 1.0842 |  |  |
|  | $\mathrm{C}-\mathrm{C}$ | 1.3933 |  |  |
| TS1 | $\mathrm{C}_{1}-\mathrm{O}$ | 1.8842 | $\mathrm{O}-\mathrm{C}_{1}-\mathrm{C}_{4}$ | 100.5 |
|  | $\mathrm{C}_{1}-\mathrm{C}_{2}$ | 1.4289 | $\mathrm{H}_{1}-\mathrm{C}_{1}-\mathrm{C}_{4}$ | 164.7 |
|  | $\mathrm{C}_{2}-\mathrm{C}_{3}$ | 1.3182 |  |  |
|  | $\mathrm{C}_{3}-\mathrm{C}_{4}$ | 1.4008 |  |  |
| TS2 | $\mathrm{C}_{1}-\mathrm{O}$ | 1.2898 | $\mathrm{O}-\mathrm{C}_{1}-\mathrm{C}_{4}$ | 164.8 |
|  | $\mathrm{C}_{1}-\mathrm{C}_{2}$ | 1.4643 | $\mathrm{H}_{1}-\mathrm{C}_{1}-\mathrm{C}_{4}$ | 105.7 |
|  | $\mathrm{C}_{2}-\mathrm{C}_{3}$ | 1.3709 |  |  |
|  | $\mathrm{C}_{3}-\mathrm{C}_{4}$ | 1.4102 |  |  |
|  | $\mathrm{C}_{1}-\mathrm{H}_{1}$ | 1.6467 |  |  |
| TS3 | $\mathrm{C}_{1}-\mathrm{O}$ | 1.2632 | $\mathrm{O}-\mathrm{C}_{1}-\mathrm{C}_{2}$ | 122.6 |
|  | $\mathrm{C}_{1}-\mathrm{C}_{2}$ | 1.5530 | $\mathrm{C}_{1}-\mathrm{C}_{5}-\mathrm{C}_{9}$ | 100.9 |
|  | $\mathrm{C}_{2}-\mathrm{C}_{3}$ | 1.4979 | $\mathrm{C}_{1}-\mathrm{C}_{5}-\mathrm{H}_{10}$ | 145.0 |
|  | $\mathrm{C}_{3}-\mathrm{C}_{4}$ | 1.4031 |  |  |
|  | $\mathrm{C}_{4}-\mathrm{C}_{5}$ | 1.3783 |  |  |
|  | $\mathrm{C}_{5}-\mathrm{C}_{6}$ | 1.4876 |  |  |
|  | $\mathrm{C}_{6}-\mathrm{C}_{1}$ | 1.7822 |  |  |
|  | $\mathrm{C}_{6}-\mathrm{C}_{2}$ | 1.4788 |  |  |
| TS4 | $\mathrm{C}_{1}-\mathrm{O}$ | 1.4155 | $\mathrm{O}-\mathrm{C}_{1}-\mathrm{C}_{4}$ | 134.34 |
|  | $\mathrm{C}_{1}-\mathrm{H}_{1}$ | 1.3102 |  |  |
|  | $\mathrm{O}-\mathrm{H}_{1}$ | 1.4155 |  |  |
|  | $\mathrm{C}_{1}-\mathrm{C}_{2}$ | 1.4872 |  |  |
|  | $\mathrm{C}_{2}-\mathrm{C}_{3}$ | 1.3606 |  |  |
|  | $\mathrm{C}_{3}-\mathrm{C}_{4}$ | 1.4246 |  |  |
| TS5 | $\mathrm{C}_{1}-\mathrm{O}$ | 1.4167 |  |  |
|  | $\mathrm{C}_{6}-\mathrm{O}$ | 1.6903 |  |  |
|  | $\mathrm{C}_{1}-\mathrm{C}_{2}$ | 1.4843 |  |  |
|  | $\mathrm{C}_{2}-\mathrm{C}_{3}$ | 1.4303 |  |  |
|  | $\mathrm{C}_{3}-\mathrm{C}_{4}$ | 1.3625 |  |  |
|  | $\mathrm{C}_{4}-\mathrm{C}_{5}$ | 1.4636 |  |  |
|  | $\mathrm{C}_{5}-\mathrm{C}_{6}$ | 1.4071 |  |  |
|  | $\mathrm{C}_{6}-\mathrm{C}_{1}$ | 1.4741 |  |  |
| TS7 | $\mathrm{C}_{1}-\mathrm{O}$ | 1.3190 |  |  |
|  | $\mathrm{C}_{1}-\mathrm{H}$ | 1.7307 |  |  |
|  | $\mathrm{C}_{6}$ - H | 1.4766 |  |  |
|  | $\mathrm{C}_{1}-\mathrm{C}_{2}$ | 1.3899 |  |  |
|  | $\mathrm{C}_{2}-\mathrm{C}_{3}$ | 1.4003 |  |  |
|  | $\mathrm{C}_{3}-\mathrm{C}^{4}$ | 1.4289 |  |  |
|  | $\mathrm{C}_{4}-\mathrm{C}_{5}$ | 1.3877 |  |  |
|  | $\mathrm{C}_{5}-\mathrm{C}_{6}$ | 1.4372 |  |  |
|  | $\mathrm{C}_{6}-\mathrm{C}_{1}$ | 1.4840 |  |  |
| TS12 | $\mathrm{C}_{1}-\mathrm{O}$ | 1.4110 |  |  |
|  | $\mathrm{C}_{6}$-O | 1.4069 |  |  |
|  | $\mathrm{C}_{1}-\mathrm{C}_{2}$ | 1.4715 |  |  |
|  | $\mathrm{C}_{2}-\mathrm{C}_{3}$ | 1.4149 |  |  |
|  | $\mathrm{C}_{3}-\mathrm{C}^{4}$ | 1.3718 |  |  |
|  | $\mathrm{C}_{4}-\mathrm{C}_{5}$ | 1.4644 |  |  |
|  | $\mathrm{C}_{5}-\mathrm{C}_{6}$ | 1.3774 |  |  |
|  | $\mathrm{C}_{6}-\mathrm{C}_{1}$ | 1.8439 |  |  |

${ }^{a}$ Experimental values for benzene $r(\mathrm{C}-\mathrm{C})=1.084 \AA, r(\mathrm{C}-\mathrm{H})=$ 1.397 Å. ${ }^{67}$
state and adduct are both $\mathrm{A}^{\prime \prime}$ with the singly occupied p-orbital of the oxygen atom parallel to the benzene ring.


One can see from the calculated bond lengths shown in Table 4 that in TS 1, the $\mathrm{C}_{2}-\mathrm{C}_{3}$ bond length is nearly that of a double


TS1



TS3


TS4


TS7


TS5


TS10


TS6


TS11


TS 12


TS13

Figure 3. Calculated structures of the important transition states using B3LYP/6-311 g(d,p).
carbon-carbon bond suggesting that of the possible resonance structures shown below, R1 is the dominant structure.


Formation of phenoxy radical (Reaction 2). This step leads to the formation of a ground-state phenoxy radical and hydrogen atom. Because of the high stability of the phenoxy radical, this step is calculated to be $6.3 \mathrm{kcal} \mathrm{mol}^{-1}$ exothermic and the calculated energy barrier is $5.3 \mathrm{kcal} \mathrm{mol}^{-1}$. Our calculated overall standard enthalpy for the formation of phenoxy radical and hydrogen atom from $\mathrm{C}_{6} \mathrm{H}_{6}+\mathrm{O}\left({ }^{3} \mathrm{P}\right)\left(\Delta_{\text {react }} H^{298}\right)$ is -16.4 $\mathrm{kcal} \mathrm{mol}{ }^{-1}$ which compares well with the available experimental value of $-14.3 \mathrm{kcal} \mathrm{mol}^{-1} .{ }^{57}$ The reverse reaction of H atom with phenoxy radical to form the adduct, $\mathbf{1}$, has a calculated barrier of $11.7 \mathrm{kcal} \mathrm{mol}^{-1}$. The rate for this reaction has not been measured but our calculated barrier is higher than the measured activation energy for the addition of H atom to benzene ( $4 \mathrm{kcal} \mathrm{mol}^{-1}$ ). ${ }^{58}$

A schematic diagram of the important orbitals for this reaction are shown below. Phenoxy radical is formed by dissociation of the $\mathrm{C}_{1}-\mathrm{H}_{1}$ bond in the adduct and a rotation of the oxygen atom p-orbital so that the singly occupied orbital is perpendicular to the phenyl ring allowing for delocalization of the radical into


H
the aromatic ring. That is, there is a shift from an $\mathrm{A}^{\prime \prime}\left(C_{2_{V}}\right.$ symmetry group) in the adduct to $\mathrm{A}^{\prime}$ in the transition state and product.

One would expect this reaction to have a late transition state with a resulting high preexponential factor. As Table 4 shows, the $\mathrm{C}_{1}-\mathrm{H}_{1}$ bond length is long ( $1.6467 \AA$ ) and the $\mathrm{O}-\mathrm{C}_{1}-\mathrm{C}_{4}$ angle is large $\left(164.8^{\circ}\right)$, suggesting that the transition state is similar to the products. This results in a (low) preexponential factor because the ring vibrational modes are relatively high due to deolcalization of the unpaired electron.
Formation of formylcyclopentadiene and subsequent reactions (Reactions 3, 10, and 11). The reaction of the triplet adduct to form triplet formylcyclopentadiene, 3, is slightly exothermic ( $-3.9 \mathrm{kcal} \mathrm{mol}^{-1}$ ) and has a $6.9 \mathrm{kcal} \mathrm{mol}^{-1}$ barrier. TS 3 is highly strained due to the existence of three- and fivemembered rings. As a result, the preexponential factor for this reaction is relatively low (Table 5). The TST rate constant for this channel from the adduct is significantly smaller at all temperatures than reaction 2 as is the RRKM rate constant. This is largely due to the difference in activation energy between these two channels. However, this difference in energy ( 1.6 kcal $\mathrm{mol}^{-1}$ ) is within the uncertainty of the CBS technique being used. Thus, it is possible that reaction 3 could play a more important role in the title reaction than our predictions indicate.

The triplet state $\left(\tilde{\mathrm{a}},{ }^{3} \mathrm{~A}^{\prime}\right)$ of $\mathbf{3}$ is a formed from the singlet ground state ( $\tilde{\mathrm{X}},{ }^{1} \mathrm{~A}^{\prime}$ ) by promoting an electron from a nonbonding p-orbital on the oxygen atom to a $\pi^{*}$-orbital on the ring as seen below. Further reaction to form cyclopentadienyl radical and HCO must be accompanied by a transition of a $\pi$ electron from the ring back to the nonbonding orbital on the O atom. For this reason, there is a slight barrier ( $8.4 \mathrm{kcal} \mathrm{mol}^{-1}$ ) to the reverse of reaction 10 and TS10 relative to $\mathbf{3}$ has an energy ( $14.4 \mathrm{kcal} \mathrm{mol}^{-1}$ ) greater than the bond dissociation energy of $3\left(6.0 \mathrm{kcal} \mathrm{mol}^{-1}\right)$. The standard enthalpy of formation of $\mathrm{C}_{5} \mathrm{H}_{5}$ +HCO is calculated to be $-7.9 \mathrm{kcal} \mathrm{mol}^{-1}$ relative to the reactants compared with and experimental value of -11.5 kcal $\mathrm{mol}^{-1} .{ }^{59}$ However, there is substantial reported uncertainty in this experimental value.


Formation of Phenol (Reaction 4). The reaction of the triplet adduct to form triplet phenol, though exothermic ( -7.9 kcal $\mathrm{mol}^{-1}$ ) has a high barrier ( $24.0 \mathrm{kcal} \mathrm{mol}{ }^{-1}$ ) typical of 1,2 hydrogen shifts ${ }^{60}$ and similar to the barrier in reaction 7. Our CBS calculations can be calibrated by comparing the calculated enthalpy of formation of singlet phenol relative to the reactants to the experimental measurements. Our calculated standard enthalpy of reaction is $-103.5 \mathrm{kcal} \mathrm{mol}^{-1}$ and the experimental value is $-102.4 \mathrm{kcal} \mathrm{mol}^{-1} .{ }^{61}$ In addition, our calculated singlet-triplet splitting in this molecule, $84.2 \mathrm{kcal} \mathrm{mol}^{-1}$, agrees with the experimental value of $81.6 \mathrm{kcal} \mathrm{mol}{ }^{-1} .{ }^{62}$ Surprisingly, the preexponential factor is high for this reaction (Table 5), due to low frequency vibrational modes in the ring. This is because the transition state is early with much adduct character (Table 5).

As an aside, because singlet phenol is the lowest singlet product from this reaction, if a transition to the singlet surface were to occur it is likely that this species would be formed.


Figure 4. Illustration of the potential energy of triplet products and transition states in the benzene and oxygen atom reaction. The transition states for products $\boldsymbol{6}, \boldsymbol{8}$, and $\boldsymbol{9}$ are off the scale of the graph (see Table 3).

TABLE 5: Transisition State Theory High Pressure Limit Calculations of Rate Constants

| reaction | $\Delta E^{298} \mathrm{TS} \mathrm{kcal} \mathrm{mol}{ }^{-1 a}$ | $\Delta S^{298} \mathrm{TS} \mathrm{cal} \mathrm{mol}{ }^{-1 b}$ | in units of $\mathrm{cm}^{3} \mathrm{~mol} \mathrm{~s}$ |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: |
|  |  |  | $A^{298} \mathrm{c}$ | $k^{298 d}$ | $k^{1800 d}$ |
| 1 forward | 4.3 | -21.7 | $2.04 \times 10^{13}$ | $1.44 \times 10^{10}$ | $6.14 \times 10^{12}$ |
| 1 reverse | 13.7 | -1.0 | $1.01 \times 10^{13}$ | $8.43 \times 10^{2}$ | $2.17 \times 10^{11}$ |
| 2 forward | 5.3 | -2.2 | $5.47 \times 10^{12}$ | $7.27 \times 10^{8}$ | $1.25 \times 10^{12}$ |
| 3 forward | 6.9 | -2.3 | $5.39 \times 10^{12}$ | $4.53 \times 10^{7}$ | $7.79 \times 10^{11}$ |
| 3 reverse | 10.1 | -7.3 | $4.31 \times 10^{11}$ | $1.67 \times 10^{4}$ | $2.56 \times 10^{10}$ |
| 4 forward | 24.2 | -0.7 | $1.22 \times 10^{13}$ | $2.11 \times 10^{-5}$ | $1.39 \times 10^{10}$ |
| 5 forward | 8.9 | -2.5 | $4.84 \times 10^{12}$ | $1.45 \times 10^{6}$ | $4.03 \times 10^{11}$ |
| 5 reverse | 1.6 | -1.6 | $7.38 \times 10^{12}$ | $4.77 \times 10^{11}$ | $4.69 \times 10^{12}$ |
| 6 forward | 33.7 | -5.7 | $9.60 \times 10^{11}$ | $1.85 \times 10^{-13}$ | $7.7 \times 10^{7}$ |
| 7 forward | 28.0 | -1.8 | $6.68 \times 10^{12}$ | $1.92 \times 10^{-8}$ | $2.65 \times 10^{9}$ |
| 7 reverse | 51.0 | -4.5 | $1.76 \times 10^{12}$ | $6.47 \times 10^{-26}$ | $1.12 \times 10^{6}$ |
| 10 forward | 14.2 | -2.0 | $6.08 \times 10^{12}$ | $2.43 \times 10^{2}$ | $1.15 \times 10^{11}$ |
| 11forward | 21.9 | -3.8 | $2.45 \times 10^{12}$ | $2.02 \times 10^{-4}$ | $5.33 \times 10^{9}$ |
| 12 forward | 18.6 | -1.1 | $9.75 \times 10^{12}$ | $2.28 \times 10^{-1}$ | $5.40 \times 10^{10}$ |
| 13 forward | 24.4 | -1.2 | $9.24 \times 10^{12}$ | $1.26 \times 10^{-5}$ | $1.02 \times 10^{10}$ |

${ }^{a}$ Energy difference between reactants and transition state (including ZPE) from CBS-QB3 calculations. ${ }^{b}$ Entropy difference between reactants and transition state, from B3LYP/6-311 g(d,p) calculations. ${ }^{c}$ Arrhenius preexponential factor determined by eq 9 . ${ }^{d}$ Rate constants determined using $\Delta E$ as Arrhenius activation energy.

The high level of chemical activation ( $-102.4 \mathrm{kcal} \mathrm{mol}^{-1}$ ) would be enough to overcome the barrier to form phenoxy radical $\left(88.9 \mathrm{kcal} \mathrm{mol}^{-1}\right)^{63}$ and one might expect that phenoxy radical formation would dominate the singlet surface.

Formation of Benzene Oxide and Oxepin (Reactions 5 and 12). The formation of benzene oxide, $\mathbf{5}$, from the adduct is exothermic ( $7.1 \mathrm{kcal} \mathrm{mol}^{-1}$ ) with a $8.9 \mathrm{kcal} \mathrm{mol}^{-1}$ barrier. The transition state for this step is also tight, leading to a low preexponential factor. As a result, this channel does not play a significant role in the title reaction. We did not find a direct route for the formation of oxepin, $\mathbf{1 3}$ from the adduct $\mathbf{1}$, although it can be formed in the triplet state from the benzene oxide, reaction 12. This step is exothermic $\left(-23.4 \mathrm{kcal} \mathrm{mol}^{-1}\right)$ with a
large barrier ( $18.6 \mathrm{kcal} \mathrm{mol}^{-1}$ ). The photochemistry of benzene oxide and oxepin has been well studied in solution ${ }^{64}$ and the gas phase ${ }^{65}$ Reactions with triplet sensitizers show that $\mathbf{1}$ can be formed from triplet benzene oxide which is consistent with our reaction barrier of $1.9 \mathrm{kcal} \mathrm{mol}^{-1}$ for this reaction. Triplet sensitized photochemsitry of benzene oxide also showed the formation of phenol. This could be due to photochemistry of the adduct or hydrogen atom abstraction by phenoxy radical. However, because TS5 is significantly higher in energy than TS2 and TS3, these reactions are unlikely to play an important role in combustion chemistry.

Formation of bicyclo[2,2,1]hepta-7-oxa-2,5-diene (Reaction 6). Our CBS calculations suggest that this reaction step is highly


Figure 5. Calculated RRKM rate constants for formation of products from the reaction of oxygen atom $\left({ }^{3} \mathrm{P}\right)$ with benzene. Experimental data ${ }^{6}$ $\left(\mathrm{O}+\mathrm{C}_{6} \mathrm{H}_{6}\right.$ _ Products) along with the sum of all reaction channels is plotted at a pressure of 100 Torr nitrogen.
endothermic, $22.2 \mathrm{kcal} \mathrm{mol}^{-1}$. Furthermore, the reaction barrier is high, $34.2 \mathrm{kcal} \mathrm{mol}{ }^{-1}$, and the preexponential is low. As a result, it is unlikely that this pathway will be an important exit channel for the chemically activated adduct.

Formation of 2,4-cyclohexadienone and n-butadienylketene (Reactions 7 and 13). The reaction of the adduct to form 2,4cyclohexadienone 7 is calculated to be exothermic ( -32.2 kcal $\mathrm{mol}^{-1}$ ), but because this reaction step is a 1,2 hydrogen shift, the barrier is high ( $28.0 \mathrm{kcal} \mathrm{mol}^{-1}$ ). Because of its high reaction barrier, this reaction pathway does not contribute significantly in combustion. The further reaction of triplet 2,4-cyclohexadienone to yield triplet $n$-butadienylketene, 13, is endothermic $21.0 \mathrm{kcal} \mathrm{mol}^{-1}$, with a calculated barrier of $24.4 \mathrm{kcal} \mathrm{mol}^{-1}$. The photochemical formation of $n$-butadienylketene from 2,4cyclohexadienone as been reported in the literature, ${ }^{64}$ but it is not clear if this reaction is on the triplet surface.

Formation of 2,5-cyclohexadienone and 2,4,5-cyclohexatrienol (Reactions 8 and 9). Reaction 8 is exothermic ( $-6.3 \mathrm{kcal} \mathrm{mol}^{-1}$ ) but the high apparent reaction barrier (36.7 kcal $\mathrm{mol}^{-1}$ ) make this reaction unlikely. Similarly, though Reaction 9 is only slightly endothermic ( $2.7 \mathrm{kcal} \mathrm{mol}^{-1}$ ), the apparent ( $62.8 \mathrm{kcal} \mathrm{mol}^{-1}$ ) barrier prevents this channel from contributing significantly. The long distances that the reactive hydrogens must travel in these reactions likely accounts for the high barriers.

Figure 5 displays the rate constants obtained from an RRKM modeling of the title reaction at 100 Torr, the pressure of the experimental values from ref 6 . The dominant product channel at higher temperatures is phenoxy plus hydrogen atom, followed closely by the elimination reaction leading to the formyl and cyclopentadienyl radicals. Formation of the stabilized formylcyclopentadiene $\mathbf{3}$ is high at lower temperatures, but as temperature increases the collisional deactivation of an excited formylcyclopentadiene molecule is less likely. The experimental values found by Nicovitch ${ }^{6}$ are in close agreement with our theoretical prediction.

## Discussion

The computational results shown above present a picture that is qualitatively and quantitatively consistent with the data generated experimentally, although there are some differences in energetics which can probably be attributed to uncertainties in the molecular modeling process. However, the interpretation of the data presented by this work is somewhat different than that which is present in the literature. Of the experimental studies, only the study of Sibener et al. ${ }^{5}$ directly observed the masses of reaction products although no study could analyze the molecular structure of the products. Sibener et al. concluded that the two primary channels for the title reaction were phenol and phenoxy +H with a minor channel being $\mathrm{CO}+$ cyclopentadiene. Our results suggest that the elimination of hydrogen directly from the triplet adduct has low barrier and is the dominant reaction at higher temperatures. Another route for hydrogen elimination does exist by first isomerizing the ring structure into formylcyclopentadiene and eliminating the tertiary hydrogen on the ring. Not only is this bond the weakest $\mathrm{C}-\mathrm{H}$ bond of the molecule, but the products are stabilized by the resonance in the cyclopentadienyl system. The resulting barrier of this elimination step is $9.5 \mathrm{kcal} \mathrm{mol}^{-1}$ relative to stable benzene and oxygen.

Sibener et al. assigned the addition product to phenol, however the flame data of Bittner and Howard ${ }^{66}$ indicate that a direct route to phenol is not possible. This work shows, instead, that the addition product is likely to be formylcyclopentadiene. The formylcyclopentadiene will be formed with a large amount of internal energy, so an elimination to form cyclopentadienyl and formyl radical is likely, especially at higher temperatures, when collisional stabilization occurs less frequently.

The observation of Sibener et al. that a minor $\mathrm{C}_{5} \mathrm{H}_{6}+\mathrm{CO}$ channel existed was based on an analysis of the angle and velocity of the scattered mass 65 signal. An alternate interpretation of the crossed molecular beam data, and more likely from our results, is that a formyl radical was produced from a chemically activated formylcyclopentadiene molecule. Thus, all the results of the molecular beam study are consistent, in a qualitative sense, with the computational work presented here. The relative distributions of the three channels cannot be directly compared with the Sibener et al. result because the molecular beams do not have thermal energy distributions.

## Conclusion

Major discrepancies between benzene flame data and elementary reaction modeling have led us to theoretically analyze the reaction pathways of the reaction $\mathrm{C} 6 \mathrm{H} 6+\mathrm{O}(3 \mathrm{P})$ using ab initio methods at the B3LYP/cc-pVDZ and CBS-QB3 level of theory. Rate constants were found from RRKM theory using the vibration frequencies and energies calculated by Gaussian 98. The overall rate constant is in close agreement to overall rate constants in the literature. The major pathway contributions at combustion temperatures are the elimination of hydrogen to form a phenoxy radical, and the formation of the cyclopentadiene and formyl radicals. Our original hypothesis that the production rate of phenoxy in our modeling of the Bittner and Howard data is too high is not supported in the present work. That is, using the rate constants derived from this study still leads to a large over-prediction in the phenoxy profile of a benzene flame. Barring errors in the current calculations, we are forced to conclude that the second hypothesis, the destruction rates and/or the destruction channels for phenoxy are in error, is correct. However, it must be noted that the uncertainty in the CBS-QB3 method is approximately $2.8 \mathrm{kcal} \mathrm{mol-1}$. If the
relative energy levels of TS2 and TS3 shifted by amounts within the error limits, the relative importance of the phenoxy and cyclopentadienyl channels would shift dramatically. Clearly, additional work is needed on phenoxy destruction channels in order test our second hypothesis.

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